 Investigation of Input-Output Waveform Engineered Continuous Inverse Class F Power Amplifiers

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Abstract—An in-depth analysis of the continuous inverse Class F power amplifier (PA) accounting for nonlinear input and output active device properties is presented. The analyses show possible ways of exploiting input nonlinearity to improve and maintain PA performance in a broadband operation and propose a flexible source second harmonic design space which reduces the input matching network (MN) design complexities. Such exploitation of input nonlinearity can also alleviate performance degradation due to dynamic knee behaviour of a practical field effect transistor (FET) in continuous inverse Class F PA operation. The analyses are validated with vector load pull (VLP) measurements and utilized to implement a broadband PA design. High drain efficiency over 75% and output power more than 38 dBm are achieved over 0.8-1.4 GHz at constant 3 dB gain compression.

Index Terms—continuous mode PA, harmonic tuned PA, input harmonics, inverse class F, multi-harmonic load pull, RF power amplifier, pulsed wave load pull, waveform engineering.

I. INTRODUCTION

HIGH EFFICIENCY power amplifier (PA) is an integral part of efficient radio frequency (RF) transmitting system design. To enable modern and future wireless communication systems, the quest for efficient, wideband, and linear PA design techniques is ongoing. Pushed by such high interest, the efficient PA design techniques are shifted from the tuned load (TL) to harmonic tuned (HT) strategies [1] leaving the job of improving linearity by digital signal processing. In the HT PAs [2]–[9], load harmonic voltage and current components are utilized to shape the drain voltage and current waveforms such that their overlapping is reduced, and device power dissipation is decreased. This is typically realized by terminating the second and third harmonic loads at either short or open (e.g. Class F, inverse Class F). Although such constraints in harmonic termination increase the matching network (MN) realization complexity they ensure an increase in PA efficiency performance [10]. Nevertheless, these design approaches are mainly for single frequency applications due to the constraints of single and fixed harmonic terminations. To overcome this and to realize wideband operation, continuous mode PA design techniques have been proposed [11]–[29]. In the continuous mode of PA operation, load harmonic terminations are swept near open or short which allows flexibility in output matching network (MN) design and realization for wideband operation and maintains PA performance over the band.

The above-mentioned works consider input source harmonic to be short circuited. Such consideration is under the fact that the impact of source harmonic terminations other than short circuit generates input nonlinearity at the intrinsic gate node and can be detrimental to PA performance [30], [31]. Thus, the input nonlinearity of the active device is avoided either by short circuit termination [32], [33] or by counteracting nonlinearity [34]. In practice, the source harmonics are designed near short to allow flexibility in MN design, but, the potential of input nonlinearity and its impact on waveform shaping has not been fully explored nor analyzed for continuous-mode HT PAs. Although, the impact of input nonlinearity is investigated in [35]–[37], a complete analysis and explanation is still missing for continuous-mode HT PAs.

In this work, a comprehensive theoretical analysis of input and output waveform engineered continuous-mode inverse Class F PA is presented. This paper expands on the related conference paper [38], which demonstrated the performance degradation of output-engineered continuous-mode inverse class F PAs due to variable knee voltage effect. Here, we present more generic analysis under input nonlinearity (both magnitude and phase) and perform a detailed investigation of the dynamic knee voltage behaviour of continuous-mode inverse class PAs in the presence of input nonlinearity.

A functional block diagram of the PA is shown in Fig. 1. In this paper, we leverage the combined benefits of input and
output harmonic manipulation and develop a continuous-mode inverse Class F PA with high efficiency over a wide frequency range. A primary contribution of this work is in demonstrating that input harmonic tuning can help to alleviate the problem of variable knee voltage behaviour of typical continuous-mode inverse Class F PAs. Additionally, we demonstrate that input nonlinearity influences the drain current waveforms, consequently, the load design space. Following the analysis and load pull validation, a wideband PA is designed and implemented covering 0.8-1.4 GHz with small signal gain more than 12 dB, efficiency over 75%, and output power more than 38 dBm at constant 3 dB gain compression.

The remainder of this paper is organized as follows. In Section II, the generalized theory of inverse Class F PA and its continuous mode of operation with input and output nonlinearity is presented. The load pull validation of the theory is presented in Section III. In Section IV, PA implementation and measurement results are reported.

II. INPUT-OUTPUT WAVEFORM ENGINEERING THEORY

The gate voltage and drain current waveforms of a tuned load PA under Class B bias condition can be expressed as [1]

\[ v_{GS,\beta}(\theta) = V_{GSO} + V_i \cos \theta \]  \hspace{1cm} (1)

\[ i_{DS,\beta}(\theta) = I_{DC,\beta} + \sum_{n=1}^{\infty} I_n \cos(n \theta) \]  \hspace{1cm} (2)

where \( V_{GSO} \) is the gate bias voltage, \( V_i \) is the fundamental gate voltage component, \( I_{max} \) is the maximum current limit of the device, and \( I_{DC,\beta} \) is the DC drain current component, and \( I_n \) denotes the DC and Fourier current components, respectively. The harmonic terminations for both input and output are considered short in (1) and (2). With second harmonic input nonlinearity at the gate node, the voltage waveform can be expressed and normalized as

\[ v_{GS}(\gamma) = v_{GS,\beta}(\theta) + v_1 \cos \theta + v_2 \cos(2\theta + \varphi_2) \]  \hspace{1cm} (3)

where \( v_2 \) is the second harmonic gate voltage component, \( \varphi_2 \) is the phase difference between fundamental and second harmonic voltage component and \( \gamma = V_2 / V_i \) is defined as the input nonlinearity factor. The gate voltage waveforms for different input nonlinearity factor, \( \gamma \), and phase difference, \( \varphi_2 \), are shown in Fig. 2.(a-d). Based on the \( \gamma \) and \( \varphi_2 \), the gate voltage waveforms are shaped and consequently, alters the drain current waveforms. The drain current waveform can be deduced considering a constant trans-conductance device for a Class B bias condition as

\[ i_{DS}(\theta, \gamma, \varphi_2) = \left\{ \begin{array}{ll}
I_{max} [\cos \theta + \gamma \cos(2\theta + \varphi_2)], & -\frac{\beta}{2} \leq \theta \leq \frac{\beta}{2} \\
0, & -\pi \leq \theta \leq -\frac{\beta}{2}, \frac{\beta}{2} \leq \theta \leq \pi
\end{array} \right. \]  \hspace{1cm} (4)

where, \( \beta \) is the modified conduction angle by gate voltage waveform shaping due to second order input nonlinearity, \( I_{DC} \) is the DC drain current component, and \( I_n \) is the Fourier drain current components of real and reactive components of \( i_{DS}(\theta) \), respectively. The effective conduction angle due to input waveform shaping can be deduced by setting \( i_{DS}(\theta) = 0 \) and can be related to \( \gamma \) and \( \varphi_2 \) as

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![Diagram of an inverse Class F PA with input nonlinearity](https://via.placeholder.com/150)---

Fig. 1. Functional block diagram of an inverse Class F PA with input nonlinearity.
The variation of the conduction angle of the device under the presence of input nonlinearity is shown in Fig. 3(a) for a Class B bias condition. The conduction angle changes widely due to input nonlinearity from ~130° to ~230° compared to the original conduction angle of 180°. Although, the conduction angle is reported to be increased by input nonlinearity (where the phase difference $\phi_2$ is considered to be 180°) in previous works [30], [39] and thus detrimental to PA efficiency, this analysis shows that the conduction angle depends on the phase difference between input fundamental and harmonic voltage components. Thus, the previous works can be considered as special cases of this work. The resultant current waveforms for different $\gamma$ and $\phi_2$ are illustrated in Fig. 2(e-h) for Class B bias condition. It is worth mentioning here that the input nonlinearity results in not only a wide variation in conduction angle, but also changes the shape of the current waveform from its original shape of half sinusoidal for a TL PA with Class B bias condition. Due to the presence of input non-linearity, current peaking, as well as the asymmetry of the drain current with respect to Y-axis, occurs which results in quadrature current components in (5). The variation of efficiency of a Class B PA with input nonlinearity is calculated and shown in Fig. 3(b). The input nonlinearity significantly impacts the efficiency performance of a Class B PA. While an optimum combination of $\gamma$ and $\phi_2$ can result in maximum efficiency of 88%, it can be degraded to the lowest value of 50% if the worst combination of $\gamma$ and $\phi_2$ is considered. When $\gamma = 0$, the efficiency in Fig. 3(b) converges to the conventional Class B PA with input and output harmonic terminations as short-circuited and shows the efficiency of 78.5%. This ensures the accuracy of the theoretical calculation.

$$\cos \frac{\beta}{2} + \gamma \cos(\beta + \phi_2) = 0$$

(6)
A. Input-Output Engineered Inverse Class F PA Analysis

For an inverse Class F PA, the second harmonic load is terminated as an open circuit at the intrinsic drain node. Although the active device generates the second harmonic drain current, it cannot be sustained due to an open circuit termination. Thus, the intrinsic drain current waveform of an inverse Class F PA is derived modifying (5) using current trapping technique as

\[
i_{DS,F-1}(\theta, \gamma, \varphi_2) = \begin{cases} 
I_{\text{max}} [\cos \theta + \gamma \cos(2 \theta + \varphi_2)] - & \frac{\beta}{2} < \theta < \frac{\beta}{2} \\
0, & -\pi < \theta < -\frac{\beta}{2}, -\frac{\beta}{2} < \theta < \pi 
\end{cases}
\]

(7)

\[
i_{DS,F-1}(\theta, \gamma, \varphi_2) = I_{\text{dc}} + \sum_{n=1}^{2m} I_{n,\varphi} \cos(n \theta) + \sum_{n=1}^{2m} I_{n,\varphi} \sin(n \theta)
\]

where, \(I_2 = \chi \cos 2\theta - \xi \sin 2\theta\), and \(\chi, \xi\) denote the coefficients for real and reactive terms of second harmonic drain current, respectively. The exclusion of second harmonic drain current components impacts the shape as well as the peak of the drain current waveform. The coefficients \(\chi\) and \(\xi\) can be computed as a function of \(\beta, \gamma,\) and \(\varphi_2\) by equating the real \(I_{2R}\) and imaginary \(I_{2I}\) part of the second harmonic current to zero as

\[
I_{2R} = \frac{1}{\pi} \int_{-\beta/2}^{\beta/2} i_{DS,F-1}(\theta) \cos 2\theta \ d\theta = 0
\]

(8)

\[
I_{2I} = \frac{1}{\pi} \int_{-\beta/2}^{\beta/2} i_{DS,F-1}(\theta) \sin 2\theta \ d\theta = 0
\]

(9)

From (6)-(9), \(\chi\) and \(\xi\) are calculated as functions of \(\beta, \gamma,\) and \(\varphi_2\) as

\[
\chi(\beta, \gamma, \varphi_2) = I_{\text{max}} \frac{3\gamma \cos \varphi_2 (2\beta + \sin 2\beta) - 16 \sin \frac{\beta}{2} + 24 \sin \frac{\beta}{2}}{6\beta + 3 \sin 2\beta}
\]

(10)

\[
\xi(\gamma, \varphi_2) = \gamma I_{\text{max}} \sin \varphi_2
\]

(11)

It is no surprise that the coefficient of reactive term, \(\xi\), is not zero unlike inverse Class F amplifier with second harmonic source impedance short-circuited, but a function of \(\gamma,\) and \(\varphi_2\) which appears due to the input second harmonic nonlinearity. The dc \(I_{\text{dc}}\), fundamental \(I_{1R}\), second harmonic \(I_{2R}\) and third harmonic \(I_{3R}\) components of drain current \(i_{DS,F-1}\) can be calculated as functions \(\beta, \gamma,\) and \(\varphi_2\) as

\[
I_{DC}(\beta, \gamma, \varphi_2) = \frac{1}{2\pi} \left( 2I_{\text{max}} \sin \frac{\beta}{2} - (\chi - \gamma I_{\text{max}} \cos \varphi_2) \sin \beta \right)
\]

(12)
\[ I_1(\beta, \gamma, \varphi_2) = \frac{I_{\text{max}}}{2\pi} \left( \beta + \sin \beta \right) - \frac{\chi}{\pi} \left( \sin \frac{\beta}{2} + \frac{1}{3} \sin \frac{3\beta}{2} \right) + \frac{\gamma I_{\text{max}} \cos \varphi_2}{3\pi} \left( 3\sin \frac{\beta}{2} + \sin \frac{3\beta}{2} \right) \]

\[ I_2(\beta, \gamma, \varphi_2) = \frac{1}{3\pi} \left[ 4\xi \sin^3 \frac{\beta}{2} - \gamma I_{\text{max}} \sin \varphi_2 \left( 3\sin \frac{\beta}{2} - \sin \frac{3\beta}{2} \right) \right] \]

\[ I_{2q}(\beta, \gamma, \varphi_2) = \frac{1}{\pi} \left[ (\xi - \gamma I_{\text{max}} \sin \varphi_2) \left( 2\beta - \sin 2\beta \right) \right] \]

\[ I_{2r}(\beta, \gamma, \varphi_2) = \frac{1}{\pi} \left[ (\xi - \gamma I_{\text{max}} \sin \varphi_2) \left( \frac{\beta}{2} + \frac{1}{3} \sin 2\beta \right) \right] \]

\[ I_{3q}(\beta, \gamma, \varphi_2) = \frac{1}{\pi} \left[ (\xi - \gamma I_{\text{max}} \sin \varphi_2) \left( \frac{\beta}{2} - \frac{1}{5} \sin 5\beta \right) \right] \]

For the higher order harmonic current components, \( n > 3 \), the coefficients can be calculated as

\[ I_{nq}(\beta, \gamma, \varphi_2) = \frac{1}{\pi} I_{\text{max}} \frac{\pi}{n-2} \left( \sin \frac{(n-1)\beta}{2} + \sin \frac{(n+1)\beta}{2} \right) - \frac{\chi}{\pi(n-2)} \left( \sin \frac{(n-2)\beta}{2} + \sin \frac{(n+2)\beta}{2} \right) - \frac{2\gamma I_{\text{max}} \cos \varphi_2}{\pi(n^2-4)} \left( \sin \beta \cos \frac{n\beta}{2} - \sin \frac{(n-2)\beta}{2} \right) \]

\[ I_{nr}(\beta, \gamma, \varphi_2) = \frac{\xi}{\pi} \frac{\pi}{n-2} \left( \sin \frac{(n-2)\beta}{2} - \sin \frac{(n+2)\beta}{2} \right) - \frac{2\gamma I_{\text{max}} \sin \varphi_2}{\pi(n^2-4)} \left( 2\cos \beta \sin \frac{n\beta}{2} - n \sin \beta \cos \frac{2\beta}{2} \right) \]

Due to the input nonlinearity, second harmonic drain current component is generated by the active device. However, that contribution is reduced out by the second harmonic open circuit loading condition for inverse Class F mode of operation. The variations of cancelling factors, \( \chi \) and \( \xi \), with input nonlinearity are shown in Fig. 4(a) and Fig. 4(b). The variation of \( \chi \) and \( \xi \) is such that the second harmonic drain current components, \( I_{2q} \) and \( I_{2r} \), are zero for inverse Class F mode of operation as shown in Fig. 4(e). Since, the DC, fundamental, and third harmonic...
Fourier coefficients are dependent on χ and ξ, the input nonlinearity also impacts their values. The variation is illustrated in Fig. 4(c), 4(d) and 4(f). By using up to first three harmonic components, the drain current can be reconstructed and shown in Fig. 5 for different input nonlinearity factor. At φ₂ = 0, drain current waveform varies significantly as shown in Fig. 5(a) due to the variation of Fourier coefficients (especially, I₇ and I₁₁) at different levels of γ. A similar trend is expected in the range of 0 < φ₂ < 90° and 270° < φ₂ < 360°. On the other hand, the coefficients are almost constant when γ changes at φ₂ = π, thus, the drain current waveforms are almost similar as shown in Fig. 5(b). Same is true in the range of 90° < φ₂ < 270°. It is worth mentioning here that the exclusion of second harmonic drain current by trapping technique also cancels out the asymmetric drain current characteristics (shown in Fig. 2) introduced by input non-linearity and make the drain current waveform symmetric with respect to Y-axis as can be seen from Fig. 5. This consequently causes reactive Fourier current coefficients (Iₙθ) of iₜₐₕ,γ,₁ to be zero.

To compute output power and efficiency, the drain voltage of inverse Class F PA is considered as [1]

$$\nu_{\text{DS,F}}(\theta) = V_{DD} - \sqrt{2}(V_{DD} - V_K) \cos \theta + \frac{1}{2\sqrt{2}} \cos 2\theta$$  \hspace{1cm} (21)

where V₃D, and Vₖ are the drain supply and device knee voltage. Thus, the DC power, output power and drain efficiency can be calculated as

$$P_{\text{DC}}(\beta, \gamma, \varphi_2) = V_{\text{DD}} \times I_{\text{DC}}(\beta, \gamma, \varphi_2)$$  \hspace{1cm} (22)

$$P_{\text{OUT}}(\beta, \gamma, \varphi_2) = \frac{1}{2} \times \sqrt{2}(V_{\text{DD}} - V_K) \times I_{\nu}(\beta, \gamma, \varphi_2)$$  \hspace{1cm} (23)

$$\eta(\beta, \gamma, \varphi_2) = \frac{1}{2} \times \sqrt{2}(V_{\text{DD}} - V_K) \times I_{\nu}(\beta, \gamma, \varphi_2)$$  \hspace{1cm} (24)

Thus, the intrinsic fundamental impedance for continuous inverse Class F PA operation as a function of second harmonic reactance, X₂, can be expressed as

$$Z_{\text{IL}} = -\frac{V_{\text{ds}} - jV_{\text{ds}}}{I_{\nu} - jI_{\nu}} = \frac{1}{Y_{\text{IL}}}$$  \hspace{1cm} (29)

where,

$$Y_{\text{IL}} = \frac{I_{\text{DC}}(\beta, \gamma, \varphi_2)}{\sqrt{2}\pi(V_{\text{DD}} - V_K)} \sigma + j \frac{8\sin^3(\beta/2)}{3\sqrt{2}X_2(2\beta - \sin 2\beta)}$$  \hspace{1cm} (30)

and,

Class B PA, the efficiency remains almost constant. Further, this analysis opens the scope of wide range of input second harmonic terminations possible instead of constant short circuit as in conventional design which is valuable in wideband PA design techniques.

B. Continuous Inverse Class F PA Analysis

For continuous-mode of inverse class F operation, the second harmonic load impedance is swept near open on the edge of the Smith chart. Thus, the second harmonic load termination can be defined as

$$Z_{\text{IL}} = -\frac{V_{\nu} - jV_{\nu}}{I_{\nu} - jI_{\nu}} = jX_2$$  \hspace{1cm} (25)

where, Vₙ and V₂ₙ are the real and reactive components of the drain voltage defined in (21), respectively, and X₂ is the second harmonic load reactance of the device. Such reactive termination at second harmonic load will allow the generation of a reactive second harmonic drain current component, I₂ₙ. From (14) and (16), Iₙ and I₂ₙ can be related as

$$I_{\nu} = I_{\nu} \frac{16}{3} \sin^2 \beta \frac{2}{2\beta - \sin 2\beta}$$  \hspace{1cm} (26)

From (25) and (26), Iₙ and I₂ₙ are defined as functions of X₂ as

$$I_{\nu} = -\frac{8}{3} \frac{V_{\text{DD}} - V_K}{X_2} \frac{\sin^2 \beta}{2\beta - \sin 2\beta}$$  \hspace{1cm} (27)

$$I_{\nu} = -\frac{1}{2} \frac{V_{\text{DD}} - V_K}{X_2}$$  \hspace{1cm} (28)

Thus, the intrinsic fundamental impedance for continuous inverse Class F PA operation as a function of second harmonic reactance, X₂, can be expressed as

$$Z_{\text{IL}} = -\frac{V_{\text{ds}} - jV_{\text{ds}}}{I_{\nu} - jI_{\nu}} = \frac{1}{Y_{\text{IL}}}$$  \hspace{1cm} (29)
From \( \gamma \), a set of load design space for continuous inverse Class F operation can be found at different values of \( \gamma \) and \( \phi \). For example, the design spaces for three different values of \( \gamma \) at \( \phi = 0, \pi/3 \), and \( \pi \) are shown in Fig. 7 where the third harmonic load termination, \( Z_{3L} \), is short-circuited. A typical trend can be observed that the fundamental load impedance is increased with a higher value of \( \gamma \) (when \( \phi < \pi/2 \)) which can be explained by the decreased amount of fundamental current component as shown in Fig. 4. Thus, the difference/separation of fundamental load design space at different \( \phi \) values is different based on the level of current component variation.

The resulting drain current waveforms for sweeping \( Z_{2L} \) near open on the edge of the Smith chart are shown in Fig. 8. Since, sweeping \( Z_{2L} \) results in drain current peaking and might exceed the device maximum limit, \( I_{\text{max}} \) (as shown in Fig. 8 for \( Z_{2L} = \pm jX_{2,\text{inv}} \)), it is important to find the maximum limit of \( X_2 \). To simplify the computation without loss of generality, the impact of \( X_2 \) variation is considered mainly on \( I_{1q} \) and \( I_{2q} \) (ignoring \( I_{3q} \) and higher order harmonics), and the drain current is written as

\[
i_{\text{DS, cont}}(\theta) = I_{\text{DC}} + I_r \cos \theta + I_{1q} \sin \theta + I_{2q} \sin 2\theta + I_{3q} \cos 3\theta
\]  

(32)

Since, the current peak does not occur at \( \theta = 0 \), the angle of maximum device current, \( \theta_{\text{max}} \), is found by equating the first derivative to zero and by confirming the second derivative at \( \theta_{\text{max}} \) to be less than zero as

\[
i_{\text{DS, cont}}'(\theta) = 0
\]  

(33a)

\[
\theta_{\text{max}} = \pm \cos^{-1}
\]

\[
\left[\begin{array}{c}
852 \cos \left(\frac{\beta}{2}\right)^2 - 1044 \cos \left(\frac{\beta}{2}\right)^4 + 400 \cos \left(\frac{\beta}{2}\right)^6 + 144 \cos \left(\frac{\beta}{2}\right)^8 - 45 \beta^2 + 360 \beta \cos \left(\frac{\beta}{2}\right)^3 \sin \left(\frac{\beta}{2}\right) - 352
\end{array}\right]
\]

\[
\sqrt{4 \cos \left(\frac{\beta}{2}\right)^8 + 10 \cos \left(\frac{\beta}{2}\right)^6 - 14 \cos \left(\frac{\beta}{2}\right)^4 + 15 \beta \sin \left(\frac{\beta}{2}\right) \cos \left(\frac{\beta}{2}\right)^3 + 12 \cos \left(\frac{\beta}{2}\right)^2 - 12}
\]

(34)
where, \( i_{DS, \text{cont}}^{*} (\theta_{\text{max}}) < 0 \) (33b)

\[
\theta_{\text{max}} = \pm \cos^{-1} \left( \frac{\sqrt{\frac{5\pi^2}{48} + \frac{22}{27}}}{2} \right) \tag{35}
\]

Similarly, \( \theta_{\text{max}} \) can be calculated for any value of \( \gamma \) and \( \varphi_2 \). Once the peak angle is determined, the limit of \( X_2 \) can be estimated by the condition \( i_{DS, \text{cont}}^{*}(\theta_{\text{max}}) \leq I_{\text{max}} \) for different value of \( \gamma \) and \( \varphi_2 \). For example, the limits of \( X_2 \) for three different conditions are derived as follows

\[
| X_2 (\gamma = 0, \varphi_2 = 0) | \geq 2.93 \times \frac{V_{DD} - V_k}{I_{\text{max}}} \tag{36a}
\]

\[
| X_2 (\gamma = 0.5, \varphi_2 = 0) | \geq 1.77 \times \frac{V_{DD} - V_k}{I_{\text{max}}} \tag{36b}
\]

\[
| X_2 (\gamma = 0.5, \varphi_2 = \pi) | \geq 2.65 \times \frac{V_{DD} - V_k}{I_{\text{max}}} \tag{36c}
\]

The drain current waveforms for the limiting value of \( X_2 = X_{2,\text{lim}} \) as defined in (36) are shown in Fig. 8. The maximum drain current reaches \( I_{\text{max}} \) at the limiting value of \( X_2 \). The variation of drain efficiency and output power in a continuous-mode of operation is shown in Fig. 9. Similar to inverse Class F PA, the efficiency performance remains almost constant with input nonlinearity compared to a Class B PA. Such behaviour allows second harmonic source termination other than a short circuit while maintains the PA efficiency performance.

C. Impact of Dynamic Knee Voltage in Continuous Inverse Class F Operation

So far, the analyses are carried out considering constant knee voltage. In reality, the device knee voltage shows a dynamic behaviour [40] with maximum device drain current. Thus, the knee voltage changes and demonstrate a dynamic behavior during the continuous inverse Class F operation due to current peaking [38]. The device knee voltage can be represented as

\[
V_k = i_{DS, \text{cont}}^{(\text{max})}, R_{\text{ON}} \tag{37}
\]

where, \( i_{DS, \text{cont}}^{(\text{max})} \) is the maximum drain current and \( R_{\text{ON}} \) is the ON resistance of the device which is mainly dependent on the mobility of the charge carriers and considered to be constant. Since the effective knee voltage of the device is increased due to current peaking of continuous inverse Class F operation, the efficiency is degraded, which is quantitatively estimated and validated with load pull measurements in [38]. However, to get a complete understanding of continuous inverse Class F PA performance, it is also important to consider the impact of dynamic knee voltage under input nonlinearity.

To compare PA performance quantitatively with and without input nonlinearity under dynamic knee voltage effect, an efficiency degradation factor \( \text{EDF} \) can be defined as

\[
\text{EDF} (\%) = 100 \times \frac{\eta(\gamma = 0, X_2 = \infty) - \eta(\gamma, \varphi_2, \beta, X_2)}{\eta(\gamma = 0, X_2 = \infty)} \tag{38}
\]

\[
\text{EDF}_{\text{max}} (\%) = 100 \times \frac{\eta(\gamma = 0, X_2 = \infty) - \eta(\gamma, \varphi_2, \beta, X_2 = X_{2,\text{lim}})}{\eta(\gamma = 0, X_2 = \infty)} \tag{39}
\]
The $E_{DF_{\text{max}}}^2$ quantifies the efficiency variation of continuous inverse Class F PA with input nonlinearity and dynamic knee voltage compared to the ideal case with second harmonic source as a short termination and no dynamic knee voltage behavior. The variation of $E_{DF_{\text{max}}}^2$ in continuous inverse Class F operation under dynamic knee voltage and input nonlinearity is shown in Fig. 10(b). With the combined effect of dynamic knee voltage and input nonlinearity, the $E_{DF_{\text{max}}}^2$ can be as high as $-10\%$. On the other hand, the $E_{DF_{\text{max}}}^2$ can also be reduced by exploiting the input nonlinearity. In fact, there is a wide range of $\gamma$ and $\phi_2$ available for which reduced $E_{DF_{\text{max}}}^2$ can be achieved. Thus, the efficiency degradation by dynamic knee voltage behaviour due to current peaking in continuous inverse Class F PA can be recovered by carefully exploiting the input nonlinearity. Based on these analyses, efficient continuous inverse Class F PA is designed in this work utilizing input nonlinearity. Such approach opens up flexible source second harmonic design space compared to the fixed short circuit condition in traditional continuous inverse Class F PA designs.

III. LOAD PULL MEASUREMENTS AND VALIDATION

A. Vector Load Pull Measurement Setup

In this work, vector load pull (VLP) measurements are performed to validate the theoretical analyses presented. VLP measurement provides useful information about device input characteristics which enhances measurement accuracy. The VLP measurement setup is shown in Fig. 11. The setup mainly consists of a load tuner MPT 1808, a source tuner MPT Lite 1808 and a phase reference unit Mesuro PR-50 from Focus Microwaves Group, VNA ZVA67 and its extension unit ZVAX-TRM40 from Rohde & Schwarz, a spectrum analyzer MS2840A from Anritsu, and DC power supplies E3634A from Agilent. The load tuner can tune the fundamental ($Z_{1L}$), second ($Z_{2L}$) and third ($Z_{3L}$) harmonic loads from 0.8 GHz to 1.8 GHz. The source tuner does the same for fundamental ($Z_{1S}$) and second ($Z_{2S}$) harmonic source impedances. To communicate, control and for synchronized measurements, Focus Device Characterization Suite (FDCS) [41] software is used. The setup mainly consists of two calibration steps. Firstly, tuner calibration and then wave calibration. The tuners are calibrated with the FDCS software using VNA ZVA67. For wave calibration, Mesuro calibration software is used along with the Mesuro phase reference unit PR-50 and a power meter NRP2 with power sensor NRP Z57 from Rohde & Schwarz. For wave measurements, both wave power and phase calibration are performed by Mesuro calibration software for $A_1$, $B_1$, $A_2$, $B_2$ as shown in Fig. 11 and utilized in the FDCS wave load-pull measurement software.

For inverse Class F PA, device compression level plays a vital role in performance metrics number. Thus, an inverse Class F PA is most often evaluated at constant compression level ensured by power sweep measurement at every load point. If continuous wave power sweeps are performed in such repetitive fashion for load pull measurements, GaN device performance are degraded due to thermal effect. To alleviate device heating and to avoid temperature effect, a pulsed wave load-pull measurement system is facilitated with a pulse modulator comes with the ZVAX-TRM40 extension unit. Such pulsed VLP measurement ensures accurate and repeatable measurements in a quasi-isothermal environment.

B. Load Pull Measurements and Validation

To validate the theoretical framework presented in the previous section, pulsed VLP measurements are performed with a Cree GaN device, CGH240010F, at 1 GHz. However, to perform accurate load pull measurements and to probe intrinsic parameters, package and device parasitics, especially the drain to source capacitance ($C_{DS}$) is needed to be extracted accurately. To do so, efficiency minima phenomenon presented in [42], [43] is utilized. Thus, the device parasitic networks along with the package parameters are extracted as shown in Fig. 12(a). A fundamental load pull measurement with inverse Class F harmonic termination as shown in Fig. 12(b) is performed at the current generator plane by de-embedding the parasitic network. A fundamental load with maximum efficiency very close to the real axis on the Smith chart as shown in Fig. 12(b) ensures the accuracy of the parasitic parameter extraction. The reason why the second and third harmonic load terminations ($Z_{2L}$ and $Z_{3L}$)

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**Fig. 11. Pulsed vector load pull measurement setup.**
are not at the perfect open and short circuit is mainly due to the loss of the coupler, fixture, and the device parasitic network.

Once the parasitic information is extracted accurately, the impact of second harmonic source pull is investigated at different fundamental and harmonic loading conditions of continuous inverse Class F operation. This is accomplished in three steps:

1. Terminate second harmonic source (\(Z_{2S}\) or \(\Gamma_{2S}\)) to short and do a fundamental load pull for different harmonic terminations of continuous inverse Class F operation. Find the maximum efficiency points.

2. Terminate fundamental load at the maximum efficiency point found in step 1 and maintain corresponding harmonic terminations.

3. Perform second harmonic source pull.

Three-second harmonic source pull results with \(Z_{2L} = \infty\), \(Z_{2L} = +j120 \, \Omega\), and \(Z_{2L} = -j120 \, \Omega\) are shown in Fig. 13(a-c), respectively. In all cases, the contour plots in Fig. 13(a-c) clearly shows that there are certain regions where efficiency degrades about 10-15% which is predicted to be \(\sim 10\%\) theoretically in Section-IIIC by \(EDF_{\text{max}}\). However, except that worst region where efficiency degrades, there is a wide range of second harmonic source terminations that show optimum efficiency performance of the device. This is unlike to the fixed short circuit termination (\(Z_{2S}\)) in conventional PA designs and can be exploited to alleviate input MN design complexities in broadband operation. However, this observation goes in line with the theoretical prediction made in Section II(A-C). The power sweep results up to constant 5 dB gain compression for different \(Z_{2S}\) points including the one provides minimum efficiency are shown in Fig. 13(d-f). As expected, the efficiency remains almost constant except the minimum efficiency point.

Intrinsic Plane
Extrinsic Plane

0.55nH 0.1nH 0.1nH

1.17pF 0.25pF 0.25pF

Fig. 12. (a) Device parasitic network, and (b) fundamental load pull at the current generator plane with inverse Class F termination.

Fig. 13. Second harmonic source pull and power sweeps for different second harmonic load termination in continuous inverse Class F operation at (a, d) \(Z_{2L} = \infty\), (b, e) \(Z_{2L} = +j120 \, \Omega\), and (c, f) \(Z_{2L} = -j120 \, \Omega\).
To investigate the impact of input nonlinearity, the input gate voltage waveforms, output drain current, and voltage waveforms are probed and closely inspected at constant 3 dB gain compression level. The waveforms are presented in Fig. 14 for different second harmonic load terminations. For inverse Class F operation, $Z_{2L} = \infty$, the source second harmonic short termination generates sinusoidal gate voltage as shown in Fig. 14(a) as expected since $\gamma = 0$. However, at other $Z_{2S}$ terminations, gate voltage waveforms deteriorate from the pure sinusoidal shape. However, the impact of input nonlinearity on the device efficiency performance depends on the combination of $\gamma$ and $\phi_2$ values as discussed in Section II. At the point of minimum efficiency, the input nonlinearity factor, $\gamma$, and the phase difference, $\phi_2$, are found to be about 0.45 and 210°. The shape of the gate voltage waveforms matches well to the one derived in Fig. 2(d). Since the input nonlinearity factor is high, the $EDF_{max}$ goes high too for a $\phi_2$ value in the range of 90° to 270° as predicted in Fig. 10(b) and causes efficiency degradation. Also, the drain current peak remains almost constant as predicted in Fig. 5 since the phase difference $\phi_2$ falls in the region of 90° to 270°. The drain voltage for all the $Z_{2S}$ terminations remains half sinusoidal as in inverse Class F operation. The input gate voltage waveform shown in Fig. 14(d) is shaped in similar fashion for $Z_{2L} = +j120 \Omega$ due to the high value of $\gamma$ about 0.65 and $\phi_2$ in the range of 180° to 270° at the minimum efficiency $Z_{2S}$ termination. For reactive $Z_{2L}$ termination, drain current peaks left as expected. On the other hand, the input nonlinearity factor, $\gamma$, and the phase difference, $\phi_2$ are found to be about 0.55 and 155° at the minimum efficiency $Z_{2S}$ termination for $Z_{2L} = -j120 \Omega$. The gate voltage waveforms are shown in Fig. 14(g). The current peaks to the right and the drain voltage waveforms remains half sinusoidal as shown in Fig. 14(h,i) for negative reactance of $Z_{2L}$ termination as expected. These results validate the theoretical analyses conducted in Section II to estimate the large signal time domain waveforms and performance variation under input-output nonlinearity in continuous inverse Class F PA operation. Most importantly, it has been observed and confirmed that there are wide ranges of $Z_{2S}$ terminations other than a short circuit possible which can maintain PA efficiency performance under input nonlinearity and dynamic knee behaviour of a practical device in continuous inverse Class F operation.

Fig. 14. Gate voltage, drain current and drain voltage waveforms for different second harmonic load termination in continuous inverse Class F operation at (a, b, c) $Z_{2L} = \infty$, (d, e, f) $Z_{2L} = +j120 \Omega$, and (g, h, i) $Z_{2L} = -j120 \Omega$. 
IV. PA DESIGN AND MEASUREMENT RESULTS

The theoretical framework and the validation with VLP measurements presented in the previous sections open a new design space for high-efficiency continuous inverse Class F PAs exploiting input nonlinearity by second harmonic source terminations other than the short. The PA design steps with input non-linearity remains same as a traditional continuous inverse Class F but with an additional harmonic source pull to select flexible clockwise Z2S terminations. This eliminates the complexities in realizing fixed short termination for broadband operation. To deploy the presented theoretical analyses, a broadband PA is designed with a Cree CG2H40010F GaN HEMT device. The fabricated prototype is shown in Fig. 15(a).

The PA is designed for the operating frequency of 0.8-1.4 GHz. The intrinsic trajectories of the source and load MN are shown in Fig. 15(b) and Fig. 15(c), respectively. The second harmonic source terminations are placed in the optimum regions as shown in Fig 13(a-c) instead of a fixed short circuit. This not only helps to maintain efficiency performance of a continuous inverse Class F PA under input-output nonlinearity and dynamic knee voltage characteristics, also reduces complexities in input MN design. A fixed second harmonic short circuit condition is no more needed.

The setup for the PA measurement is shown in Fig. 16(a). The same setup is used for both continuous and modulated signal measurements. The setup consists of a driver, isolator, a signal generator MXG N5182A from Agilent Technologies to...
Fig. 17. Spectrum of the input, output and linearized response of the PA.

This work presents a comprehensive analysis under the presence of input-output non-linearity and dynamic knee behavior of a practical FET device of a broadband inverse Class F PA design. Theoretical analysis shows new second harmonic source design space for continuous inverse Class F PA by exploiting input nonlinearity. This paper introduces a flexible source harmonic termination for broadband PA operation compared to fixed short circuit termination in traditional continuous inverse Class F PAs. Thus, the proposed PA design space reduces input MN design complexities for broadband operation. The theoretical analyses are validated with VLP measurements and the application is demonstrated with a broadband PA design exhibiting high efficiency over a wide frequency range.

REFERENCES


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